

Supersensitive, Ultrafast, and Broad-Band Light-Harvesting Scheme Employing Carbon Nanotube/TiO₂ Core–Shell Nanowire Geometry

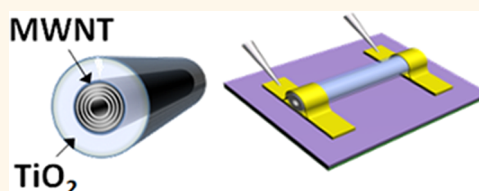
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One-dimensional metal oxide nanostructures as photodetectors (PDs) exhibiting high photoconductive gain have stimulated much interest due to the pronounced surface effects (*i.e.*, oxygen adsorption/desorption) originating from the high surface-to-volume ratio.^{1,2} The photogain as high as 10^8 has been reported using ZnO nanowires (NWs).^{3,4} In an attempt to achieve high-performance photodetection, most studies have focused on boosting the photosensitivity.^{5,6} Although the high photogain is achieved, these surface effects act as a double-edged sword that prolongs the response/recovery time due to slow oxygen adsorption/desorption processes, thus hindering practical applications of metal oxide nanostructures in PDs. Several approaches to break this compromise for enhancing the responsibility as well as reducing the response/recovery times have been reported, including an introduction of Schottky junction,⁷ the modification of surfaces,^{5,6} the reduction of nanowire sizes,⁸ and the fabrication of network structures.⁹ However, to realize the nanostructured metal oxide PDs in practical use, further improvement toward higher gain and faster response is urgently demanded.

The heterojunction between two different nanomaterials in a variety of nanostructures, such as core–shell NWs,^{10–13} segmented NWs,^{14,15} and core–shell quantum dots,¹⁶ has shown great potential for the application in solar cells, biosensors, photodetection, and piezoelectric nanogenerators.^{13,17–21} For example, the semiconducting core–shell NW geometry can enhance the efficiency of charge collection by shortening the paths traveled by minority carriers

ABSTRACT



We demonstrate a novel, feasible strategy for practical application of one-dimensional photodetectors by integrating a carbon nanotube and TiO₂ in a core–shell fashion for breaking the compromise between the photogain and the response/recovery speed. Radial Schottky barriers between carbon nanotube cores and TiO₂ shells and surface states at TiO₂ shell surface regulate electron transport and also facilitate the separation of photogenerated electrons and holes, leading to ultrahigh photogain ($G = 1.4 \times 10^4$) and the ultrashort response/recovery times (4.3/10.2 ms). Additionally, radial Schottky junction and defect band absorption broaden the detection range (UV–visible). The concept using metallic core oxide–shell geometry with radial Schottky barriers holds potential to pave a new way to realize nanostructured photodetectors for practical use.

KEYWORDS: core–shell · TiO₂ · carbon nanotube · photodetector · photoresponse time

and therefore can promote energy-harvesting properties.^{17–20} The NWs with semiconducting cores encapsulated by dielectric shells not only provide excellent surface passivation but also offer excellent light-harvesting characteristics, such as broad-band working ranges, omnidirectionality, and polarization insensitivity since the dielectric shell is capable of creating a smooth gradient of refractive index profiles.²¹

In this study, we used metallic multi-walled carbon nanotubes (MWNTs) as inner cores and semiconducting TiO₂ as outer shells to create 1D core–shell nanostructures for supersensitive, ultrafast, and

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broad-band photodetection. Among metal oxide candidates, TiO_2 , a wide band gap material (3.2 eV for anatase), is promising for the photodetection with a fast response speed and a high gain and thus could efficiently serve as a photon adsorption layer in this unique structure by enhancing charge separation promoted by the surface states.²² On the other hand, MWNT cores provide radial Schottky barriers to TiO_2 shells to regulate electron transport and also facilitates the separation of photogenerated electrons and holes. Accordingly, the photogain up to 1.4×10^4 can be achieved, and the response/recovery times are promoted to 4.3/10.2 ms. Broad-band detection is also accomplished by the additional absorption from MWNTs possibly due to the radial Schottky junction between MWNTs and TiO_2 shells and the defect band absorption. The concept of the core–shell nanowire geometry reported herein explores a new pathway toward ultrafast, supersensitive, and broad-band photodetection.

RESULTS AND DISCUSSION

Aligned MWNTs were synthesized at 1400 °C using ferrocene-dissolved xylene (1:50, 100 mL) as the carbon source.^{23,24} Then, atomic layer deposition (ALD) was employed to deposit TiO_2 on the MWNTs,²⁴ as shown in the Supporting Information. With precise control of TiO_2 growth, a 5 nm thick TiO_2 layer was deposited on MWNTs to ensure sufficient light absorption. Since the photon-to-electron conversion efficiency is highly dependent on the crystallinity of the active layer (*i.e.*, TiO_2),^{1,25} an annealing process was performed at 450 °C in Ar for 2 h. Moreover, according to the previous reports,^{26,27} 5 nm thickness could absorb about 10–15% of light from visible to UV region. Since the incident light will penetrate the TiO_2 layer twice due to the core–shell geometry, we could expect that the light absorption will be over 20%, which is sufficient for photoexcitation.

A single TiO_2 –MWNT core–shell PD was fabricated with the assistance of electron-beam lithography (EBL), as shown in a scanning electron microscopy (SEM) image and the schematic in Figure 1a (see Supporting Information for fabrication details). The core–shell NWs were dispersed in isopropyl alcohol and then transferred onto the SiO_2/Si substrate by dropping NWs from a pipet. To contact a core–shell single NW, the electrode patterns were defined by SEM and then followed by Ti/Au (30 nm/70 nm) electrode fabrication *via* the EBL process. Before testing the performance of the MWNT/ TiO_2 core–shell NW PD, the material properties were characterized. Figure 1b shows the analyses of X-ray diffraction (XRD) of core–shell NWs before and after the annealing. Before annealing, no significant feature was observed, indicating the poor crystallinity of the TiO_2 layer, while after the heat treatment, evident peaks at 24.5 and 37° were observed,

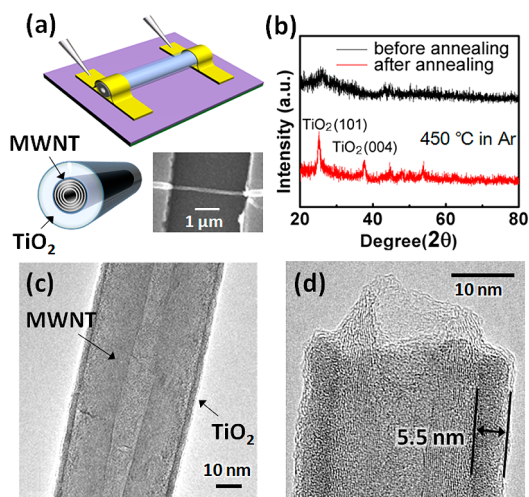


Figure 1. (a) Schematics of the core–shell nanostructure and the corresponding PD. Inset shows the SEM image of the device. (b) XRD spectra of the core–shell NW before and after the annealing process. (c) TEM image of a core–shell NW. (d) High-resolution TEM image of a core–shell NW at the end.

corresponding to (101) and (004) planes of the anatase TiO_2 , respectively. Transmission electron microscopy (TEM) further displays the features of the core–shell nanostructures. A uniformly coated TiO_2 shell is observed in Figure 1c, where the TiO_2 is darker because of heavier molecular weight, and the inner structures with stripe fringes are the walls of the MWNT. Further analysis shows that the MWNTs have diameters ranging from 20 to 60 nm, and the thickness of TiO_2 shells is ~5 nm (Figure 1d), consistent with the ALD fabrication parameters.

The electrical measurements were carried out at room temperature using a semiconductor parameter analyzer (Keithley Instruments 4200), and the photoresponse was characterized under air mass 1.5 global (AM 1.5G) illumination (1000 W/m^2) and a halogen lamp coupled to a monochromator. Before verifying the performance of the MWNT/ TiO_2 core–shell NW PD, the role of the core MWNT was first examined. A single MWNT device without TiO_2 shells, serving as a control, was made *via* the same process, and its *IV* characteristic is shown in Figure 2a. The quasi-linear *IV* curve, representing the most classic electrical behavior of MWNT devices, reveals the fact that the MWNT is metallic. As for the photoresponse of the MWNT device, the *IV* curve under AM 1.5G illumination remains the same as that in the dark, indicating that the metallic MWNT is insensitive to light.²⁸ The photoresponse of the MWNT/ TiO_2 core–shell NW is shown in Figure 2b. As compared to the pure MWNT device, the MWNT/ TiO_2 core–shell NW PD exhibits an apparent response to AM 1.5G light, demonstrating the role of TiO_2 in the photocurrent generation. Note that the contacts between the Ti electrodes with work function ϕ_m of ~4.3 eV and the TiO_2 shell with electron affinity of ~4.3 eV are ohmic.²⁹

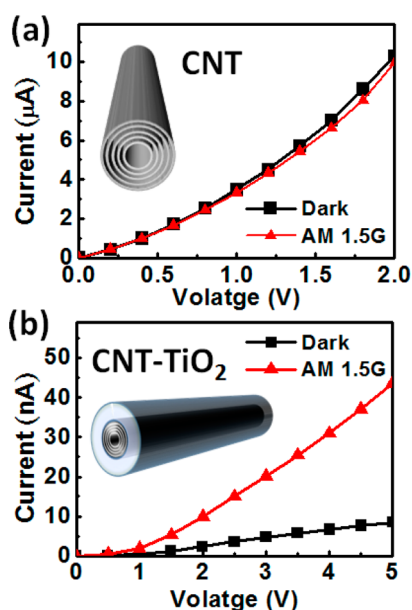


Figure 2. (a) *IV* characteristics of the MWNT device in the dark and upon AM 1.5G illumination. (b) *IV* characteristics of the TiO₂-coated MWNT PD in the dark and upon AM 1.5G illumination.

To highlight the ultrafast response of the devices, the time-resolved measurements were performed at a fixed bias of 5 V. As shown in Figure 3a, a fully reversible response was acquired; the current abruptly increased and then decreased to its initial value as the light was on and off, respectively. The amplification ratios of photocurrents range from 160 to 250%, depending on the optical wavelengths and the incident power. Note that the power density of a halogen lamp used here is varied with the wavelength. For quantitative analysis, the wavelength-dependent photogain will be carried out later. To probe the temporal response in a high resolution, we set up a data acquisition system (DAQ; DAQ-2214, ADLINK) operated with a resolution as low as 10 μ s. One can see that a short response/recovery time regardless of the wavelength is obtained, as shown in an example in Figure 3b,c at an incident wavelength of 350 nm. Here the data are partially skipped in order to more clearly reveal the difference between response and recovery speed; the interval of each data point presented in the figures is 1 ms. The response and the recovery times can be used to evaluate the speed of response and recovery processes, which are defined as the time needed to rise to $(1 - 1/e)$ of the dark current and recover to $1/e$ of the maximum photocurrent, respectively.³⁰ Accordingly, 4.3 ms of the response and 10.2 ms of the recovery times can be obtained. To the best of our knowledge, this is the fastest response/recovery speed reported for any TiO₂-based PD to date.²²

The importance of metallic core–semiconducting shell geometry in the photodetection can also be understood by the spectral photoresponse measurements, as

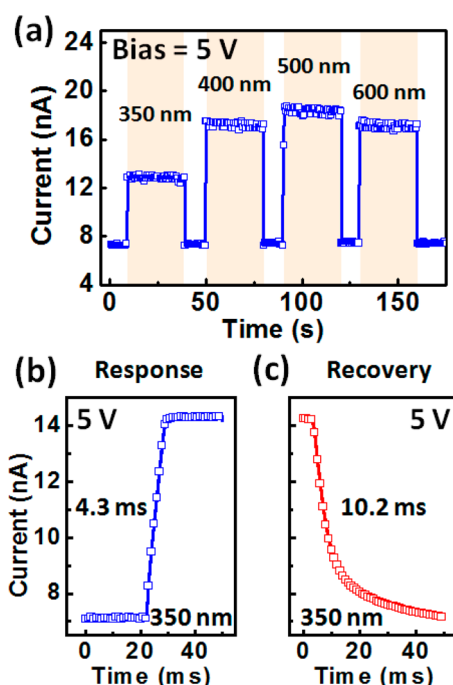


Figure 3. (a) Time-resolved photocurrents of the TiO₂-coated device under various wavelengths illumination at 5 V. (b) Highlights of the response speed under 350 nm illumination at 5 V. (c) Highlights of the recovery speed under 350 nm illumination at 5 V.

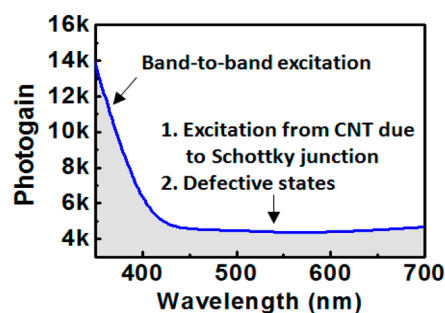


Figure 4. Photogain of the core–shell NW PD as a function of wavelengths. Three mechanisms are depicted on the graph, which can be attributed to different absorption ranges.

presented in Figure 4. The photogain is calculated by $G = (\Delta I_{ph}/q)/(P/h\nu)$, where ΔI_{ph} is the difference between measured photocurrent and dark current, q is the electron charge, P is the power of incident light absorbed by the device, and $h\nu$ is the photon energy. To obtain the power of incident light absorbed by the device, the light absorption area of the device is estimated according to the SEM and the TEM images of the device, where the length and the width of the nanostructure is 3.2 μ m and 65 nm, respectively. Note that the incident power is dependent on the incident wavelength and thus is normalized as computing the photoconductive gain (see Supporting Information). The photogain has a ramp at UV region because the TiO₂, a wide band gap semiconductor (3.2 eV), has remarkable UV absorption property. Interestingly, the

photoresponse in the regime higher than 400 nm is also obvious and could be attributed to two mechanisms associated with the structure effects: (i) Because of the radial Schottky junction between the MWNT core and the TiO₂ shell, the device possesses an exclusive light absorption pathway; that is, the photocarriers are excited from metallic MWNTs over the Schottky barrier.³¹ Therefore, in addition to band-to-band excitation in the TiO₂ shell, the electrons could be excited from the metallic MWNT to the conduction band of the TiO₂ shell, resulting in additional broad-band absorption. To undergo this process, the incident photon energy must be higher than the blocking barrier (ϕ_b)³² that originates from the radial Schottky contact between the metallic MWNT core (work function $\phi_m = 4.9$ eV)³³ and the semiconducting n-type TiO₂ shell (electron affinity $\chi = 4.3$ eV).^{34–36} According to the equation $\phi_b = \chi - \phi_m$, the ideal ϕ_b between MWNTs and TiO₂ = 0.6 eV is low, thus the absorption spectrum can be widely extended toward long wavelength region. (ii) Another possible factor for observed broad-band absorption is the presence of sub-band gap due to the defects in TiO₂ or the interfacial defects at the MWNT–TiO₂ junction.^{37–42} Those defect-induced states can broaden the absorption spectra and have been widely demonstrated in other heterojunction systems.⁴³ Note that, on the basis of current data, we cannot verify which mechanism dominates because the resultant effect might be a combination of both.

In comparison to other TiO₂-based PDs with high photogains reported elsewhere, 1.4×10^4 of the maximum photogain of the core–shell NW in the UV region (Figure 4) is about 4 orders of magnitude higher than those of TiO₂ nanotube arrays (~ 52).⁴⁴ For metal oxide NWs, it is known that the process of oxygen adsorption [$O_{2(g)} + e^- \rightarrow O_{2(ad)}^-$] and desorption [$O_{2(ad)}^- + h^+ \rightarrow O_{2(g)}$] on the NW surfaces dominates the photocarrier transport and the photoresponse behavior.^{3,4} The origin of the photogain of metal oxide nanostructures is closely connected to the long lifetime of the photogenerated electrons, resulting from the hole-trapping effect activated by the oxygen desorption at the surfaces, which significantly reduces the recombination rate of electron–hole pairs. Accordingly, the photogain can also be expressed as $G = \tau/\tau_t$, where τ is the carrier lifetime (recovery time) and τ_t is the transit time between two electrodes.⁴⁵ Since τ_t is a function only related to experimental parameters and material mobility, this equation implies a trade-off between the photogain and the response speed ($\propto 1/\tau$). This can be confirmed by the fact that most of the PDs with a high photogain exhibit a slow response speed ranging from tens of seconds to several hours.^{1,3,7} A major breakthrough in this study is that the MWNT–TiO₂ core–shell NW not only retains the high gain but also possesses high response speed, diminishing

the compromise between photogain and response speed.

Possessing ultrahigh photogains and ultrashort response/recovery times at the same time could be attributed to the unique core–shell geometry of the MWNT–TiO₂ NWs. Full depletion of ~ 5 nm thick TiO₂ is formed due to (i) the space charge layer formed at the surface of TiO₂ shells and (ii) the radial Schottky barrier between MWNT cores and TiO₂ shells, which can be confirmed by the very low dark current of the core–shell NW device, as shown in Figure 2b. Under illumination, the processes of oxygen desorption/adsorption at the surface and the radial Schottky barrier modulation effectively activate a reversible regulation of depletion width and height, leading to a great enhancement of the photogain. The rapid photocurrent response and recovery are related to the fact stated as follows. Under illumination, the strong local electric field in the TiO₂ core assisted by the oxygen desorption/adsorption and the radial Schottky junctions will quickly separate the photon-generated electrons and holes, giving rise to a rapid response speed, as shown in Figure 3b. Under illumination, the increase of the photon-generated carrier density leads to a reduction of the barrier height and a narrowing of the barrier width, enabling the access of the photo-induced low-resistive conducting path within the depleted TiO₂ and thus increasing the conductivity of the core–shell NW. It has been known that the process of the light-induced Schottky barrier modulation is rather fast.⁹ Accordingly, the response/recovery time can be dramatically reduced. We note that, in previous studies, two types of Schottky contacts pertaining to NW PDs have been demonstrated, including the Schottky contact between planar metal electrodes and NWs⁷ and the NW-to-NW contacts within NW networks,⁹ both of which have shown the enhancement to response speed. Moreover, we envision that the photoresponse of our devices can be further improved after optimizing the thickness of the TiO₂ layer for maximizing the light absorption and the depleted region. Moreover, using the aligned array of CNT/TiO₂ core–shell NWs that capitalizes on a strong light-trapping effect at a variety of wavelengths and angles of incidence might significantly increase the amount of photons absorbed at the radial junction regions due to its subwavelength structure.^{46,47} In addition, Cao *et al.* reported the use of the 1D nanostructure resonance property by tuning the radius of nanowires so that the light absorption can be enhanced at resonance regions, referred to as the leaky mode resonance enhancement.^{48,49} Implementing the above-mentioned concepts into CNT–TiO₂ core–shell PDs might be an important concern for applying it to practical applications in the future.

In summary, we break the compromise between photogains and response/recovery times of nanowire PDs using metallic MWNT–semiconducting TiO₂

core–shell nanowires. MWNT–TiO₂ NW PDs exhibit high photogains of up to 1.4×10^4 and short response/recovery times of 4.3/10.2 ms with the broad-band detection. The ultrasensitive and ultrafast photodetection achieved in MWNT–TiO₂ core–shell NWs is due to enhanced charge separation promoted by the surface states and the radial Schottky barrier between MWNT cores and TiO₂ shells. The broad-band detection results from additional absorption from MWNTs due to Schottky junction and defect band absorption. The resultant performance is essentially useful for fabricating photovoltaic and photodetection devices based on core–shell nanowire heterojunctions.

Conflict of Interest: The authors declare no competing financial interest.

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Supporting Information Available: Synthesis and device fabrication of MWNT arrays and MWNT–TiO₂ core–shell structures. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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